

## Partial differential equations

### The heat equation

We wish to describe one-dimensional heat flow.

**Comment.** If this sounds very specialized, it might help to know that the heat equation is also used, for instance, in probability (Brownian motion), financial math (Black-Scholes), or chemical processes (diffusion equation).

Let  $u(x, t)$  describe the temperature at time  $t$  at position  $x$ .

If we model a heated rod of length  $L$ , then  $x \in [0, L]$ .

**Notation.**  $u(x, t)$  depends on two variables. When taking derivatives, we will use the notations  $u_t = \frac{\partial}{\partial t}u$  and  $u_{xx} = \frac{\partial^2}{\partial x^2}u$  for first and higher derivatives.

Experience tells us that heat flows from warmer to cooler areas and has an averaging effect.

Make a sketch of some temperature profile  $u(x, t)$  for fixed  $t$ .

As  $t$  increases, we expect maxima (where  $u_{xx} < 0$ ) of that profile to flatten out (which means that  $u_t < 0$ ); similarly, minima (where  $u_{xx} > 0$ ) should go up (meaning that  $u_t > 0$ ). The simplest relationship between  $u_t$  and  $u_{xx}$  which conforms with our expectation is  $u_t = k u_{xx}$ , with  $k > 0$ .

#### (heat equation)

$$u_t = k u_{xx}$$

Note that the heat equation is a linear and homogeneous **partial differential equation**.

In particular, the principle of superposition holds: if  $u_1$  and  $u_2$  solve the heat equation, then so does  $c_1 u_1 + c_2 u_2$ .

**Higher dimensions.** In higher dimensions, the heat equation takes the form  $u_t = k(u_{xx} + u_{yy})$  or  $u_t = k(u_{xx} + u_{yy} + u_{zz})$ . The heat equation is often written as  $u_t = k \Delta u$ , where  $\Delta = \frac{\partial^2}{\partial x^2} + \frac{\partial^2}{\partial y^2} + \frac{\partial^2}{\partial z^2}$  is the Laplace operator you may know from Calculus III.

The Laplacian  $\Delta u$  is also often written as  $\Delta u = \nabla^2 u$ . The operator  $\nabla = (\partial/\partial x, \partial/\partial y)$  is pronounced “nabla” (Greek for a certain harp) or “del” (Persian for heart), and  $\nabla^2$  is short for the inner product  $\nabla \cdot \nabla$ .

Let us think about what is needed to describe a unique solution of the heat equation.

- **Initial condition** at  $t = 0$ :  $u(x, 0) = f(x)$  (IC)

This specifies an initial temperature distribution at time  $t = 0$ .

- **Boundary condition** at  $x = 0$  and  $x = L$ : (BC)

Assuming that heat only enters/exits at the boundary (think of our rod as being insulated, except possibly at the two ends), we need some condition on the temperature at the ends. For instance:

- $u(0, t) = A, u(L, t) = B$

This models a rod where one end is kept at temperature  $A$  and the other end at temperature  $B$ .

- $u_x(0, t) = u_x(L, t) = 0$

This models a rod whose ends are insulated as well.

Under such assumptions, our physical intuition suggests that there should be a unique solution.

**Important comment.** We can always transform the case  $u(0, t) = A, u(L, t) = B$  into  $u(0, t) = u(L, t) = 0$  by using the fact that  $u(t, x) = ax + b$  solves  $u_t = k u_{xx}$ . Can you spell this out?

**Example 160.** To get a feeling, let us find some solutions to  $u_t = ku_{xx}$ .

- $u(x, t) = ax + b$  is a solution.
- For instance,  $u(x, t) = e^{kt}e^x$  is a solution.  
[Not a very interesting one for modeling heat flow because it increases exponentially in time.]
- More interesting are  $u(x, t) = e^{-kt}\cos(x)$  and  $u(x, t) = e^{-kt}\sin(x)$ .
- More generally,  $e^{-k\lambda^2 t}\cos(\lambda x)$  and  $e^{-k\lambda^2 t}\sin(\lambda x)$  are solutions.
- Can you find further solutions?

**Important observation.** This reveals a strategy for solving the heat equation together with the following boundary and initial conditions:

$$\begin{aligned} u_t &= ku_{xx} && \text{(PDE)} \\ u(0, t) &= u(L, t) = 0 && \text{(BC)} \\ u(x, 0) &= f(x), \quad x \in (0, L) && \text{(IC)} \end{aligned}$$

Note that  $e^{-k\lambda^2 t}\sin(\lambda x)$  solves the PDE and also satisfies (BC) if  $\lambda = n\frac{\pi}{L}$  for some integer  $n$ . Hence,

$$u_n(x, t) = e^{-k\left(\frac{\pi n}{L}\right)^2 t} \sin\left(\frac{\pi n}{L} x\right)$$

satisfies the PDE as well as (BC) for any integer  $n$ .

It remains to satisfy (IC) and we plan to do so by taking the right combination of the  $u_n(x, t)$ . At  $t = 0$ , we get  $u_n(x, 0) = \sin\left(\frac{\pi n}{L} x\right)$  and all of these are  $2L$ -periodic and odd. This matches exactly the terms we get when we write  $f(x)$  as a Fourier sine series ( $f(x)$  is only given on  $(0, L)$  and we extend it to an odd  $2L$ -periodic function):

$$f(x) = \sum_{n \geq 1} b_n \sin\left(\frac{\pi n}{L} x\right)$$

Consequently, (PDE)+(BC)+(IC) is solved by

$$u(x, t) = \sum_{n=1}^{\infty} b_n u_n(x, t) = \sum_{n=1}^{\infty} b_n e^{-\left(\frac{\pi n}{L}\right)^2 kt} \sin\left(\frac{\pi n}{L} x\right).$$

**Comment.** Note that the coefficients  $b_n$  can be computed as

$$b_n = \frac{1}{L} \int_{-L}^L f(x) \sin\left(\frac{n\pi x}{L}\right) dx = \frac{2}{L} \int_0^L f(x) \sin\left(\frac{n\pi x}{L}\right) dx,$$

where the first integral makes reference to the extension of  $f(x)$  while the second integral only uses  $f(x)$  on its original interval of definition.

**Comment.** Note that  $n = 0$  just gives the zero function  $u_0(x, t) = 0$ , and negative values don't give anything new because  $u_{-n}(x, t) = -u_n(x, t)$ .

**Example 161.** Find the unique solution  $u(x, t)$  to:  $u_t = u_{xx}$  (PDE)  
 $u(0, t) = u(\pi, t) = 0$  (BC)  
 $u(x, 0) = \sin(2x) - 7\sin(3x), \quad x \in (0, \pi)$  (IC)

**Solution.** This is the case  $k = 1, L = \pi$  of the above. Hence, as we just observed, the functions

$$u_n(x, t) = e^{-n^2 t} \sin(nx)$$

satisfy (PDE) and (BC) for any integer  $n$ .

Since  $u_n(x, 0) = \sin(nx)$ , we have

$$u_2(x, 0) - 7u_3(x, 0) = \sin(2x) - 7\sin(3x)$$

as needed for (IC).

Therefore, (PDE)+(BC)+(IC) is solved by

$$u(x, t) = u_2(x, t) - 7u_3(x, t) = e^{-4t} \sin(2x) - 7e^{-9t} \sin(3x).$$

**Example 162.** Find the unique solution  $u(x, t)$  to:  $u_t = 3u_{xx}$  (PDE)  
 $u(0, t) = u(4, t) = 0$  (BC)  
 $u(x, 0) = 5\sin(\pi x) - \sin(3\pi x), \quad x \in (0, 4)$  (IC)

**Solution.** This is the case  $k = 3, L = 4$  of the above. Hence, the functions

$$u_n(x, t) = e^{-3\left(\frac{\pi n}{4}\right)^2 t} \sin\left(\frac{\pi n}{4} x\right)$$

satisfy (PDE) and (BC) for any integer  $n$ . Since  $u_n(x, 0) = \sin\left(\frac{\pi n}{4} x\right)$ , we have

$$5u_4(x, 0) - u_{12}(x, 0) = 5\sin(\pi x) - \sin(3\pi x),$$

which is what we need for the right-hand side of (IC). Therefore, (PDE)+(BC)+(IC) is solved by

$$u(x, t) = 5u_4(x, t) - u_{12}(x, t) = 5e^{-3\pi^2 t} \sin(\pi x) - e^{-27\pi^2 t} \sin(3\pi x).$$